ME/SE 740

Lecture 9

Euler Angles and Euler's Theorem

Euler Angles

Today we return to our examination of rotations of rigid bodies. We begin with a discussion about "Euler Angles." Consider the two rotations depicted in the figure below where the "north pole" is first rotated about the z-axis through an angle θ and this is followed by a rotation about the y-axis through an angle ϕ . The north pole ends up at point Q on the unit sphere.

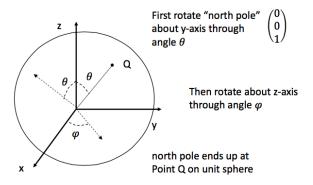


Figure 1: Two Rotations of the North Pole

These two rotations can be represented by:

$$\begin{pmatrix}
c\phi & -s\phi & 0 \\
s\phi & c\phi & 0 \\
0 & 0 & 1
\end{pmatrix}
\begin{pmatrix}
c\theta & 0 & s\theta \\
0 & 1 & 0 \\
-s\theta & 0 & c\theta
\end{pmatrix}
\begin{pmatrix}
0 \\
0 \\
1
\end{pmatrix}$$

Suppose we are given an arbitrary 3×3 rotation (of the north pole):

$$\left(\begin{array}{ccc}
n_x & o_x & a_x \\
n_y & o_y & a_y \\
n_z & o_z & a_z
\end{array}\right)
\left(\begin{array}{c}
0 \\
0 \\
1
\end{array}\right)$$

and that both map the north pole $(0,0,1)^T$ to the same point:

$$\begin{pmatrix} c\phi & -s\phi & 0 \\ s\phi & c\phi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} c\theta & 0 & s\theta \\ 0 & 1 & 0 \\ -s\theta & 0 & c\theta \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \begin{pmatrix} n_x & o_x & a_x \\ n_y & o_y & a_y \\ n_z & o_z & a_z \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

Multiplying on the left both sides of this equation by the matrix inverses (in the appropriate order) we can write:

$$\begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \underbrace{\begin{pmatrix} c\theta & 0 & -s\theta \\ 0 & 1 & 0 \\ s\theta & 0 & c\theta \end{pmatrix} \begin{pmatrix} c\phi & s\phi & 0 \\ -s\phi & c\phi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} n_x & o_x & a_x \\ n_y & o_y & a_y \\ n_z & o_z & a_z \end{pmatrix}}_{T} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

where T maps $(0,0,1)^T$ to $(0,0,1)^T$.

Given a left action of a group $G: M \longrightarrow M$ on a manifold M and a point $m \in M$ the set of elements $g \in G$ such that $g \circ m = m$ is a subgroup called the <u>isotropy</u> subgroup of M. We say that G has a left action on M if for every $g \in G$, $g_1 \circ (g_1 \circ m) = (g_1 \circ g_2)m$ for $g_1, g_2 \in G$ and $m \in M$. In particular, the isotropy subgroup of SO3) (3 × 3 proper rotation matrices) corresponding to $(0,0,1)^T$ is:

$$\left(\begin{array}{ccc} c\psi & -s\psi & 0\\ s\psi & c\psi & 0\\ 0 & 0 & 1 \end{array}\right) = \Psi$$

There is a value of ψ , with $0 \le \psi 2\pi$ such that

$$\begin{pmatrix} c\theta & 0 & -s\theta \\ 0 & 1 & 0 \\ s\theta & 0 & c\theta \end{pmatrix} \begin{pmatrix} c\phi & s\phi & 0 \\ -s\phi & c\phi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} n_x & o_x & a_x \\ n_y & o_y & a_y \\ n_z & o_z & a_z \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix} = \begin{pmatrix} c\psi & -s\psi & 0 \\ s\psi & c\psi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ 1 \end{pmatrix}$$

Solving for $\vec{n}, \vec{o}, \vec{a}$ we obtain:

$$\begin{pmatrix} n_x & o_x & a_x \\ n_y & o_y & a_y \\ n_z & o_z & a_z \end{pmatrix} = \underbrace{\begin{pmatrix} c\phi & -s\phi & 0 \\ s\phi & c\phi & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{0 \le \phi \le 2\pi} \underbrace{\begin{pmatrix} c\theta & 0 & s\theta \\ 0 & 1 & 0 \\ -s\theta & 0 & c\theta \end{pmatrix}}_{0 \le \theta \le \pi} \underbrace{\begin{pmatrix} c\psi & -s\psi & 0 \\ s\psi & c\psi & 0 \\ 0 & 0 & 1 \end{pmatrix}}_{0 \le \psi \le 2\pi}$$

Tait-Bryan Angles

$$\begin{pmatrix} n_x & o_x & a_x \\ n_y & o_y & a_y \\ n_z & o_z & a_z \end{pmatrix} = \begin{pmatrix} c\phi & -s\phi & 0 \\ s\phi & c\phi & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} c\theta & 0 & s\theta \\ 0 & 1 & 0 \\ -s\theta & 0 & c\theta \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & c\psi & -s\psi \\ 0 & s\psi & c\psi \end{pmatrix}$$

Euler's Theorem (for us). Every 3×3 proper rotation $X \in SO(3)$ is a rotation about an axis by a certain amount $\theta, 0 \le \theta \le \pi$.

Suppose that X is given:

$$X = \begin{pmatrix} \bar{n}_x & \bar{o}_x & \bar{a}_x \\ \bar{n}_y & \bar{o}_y & \bar{a}_y \\ \bar{n}_z & \bar{o}_z & \bar{a}_z \end{pmatrix}$$

By Euler's Theorem, there is a unit vector \vec{k} and a rotation $0 \le \theta \le \pi$ such that X is a rotation of θ about \vec{k} .

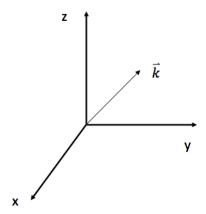


Figure 2: The axis \vec{k}

Let C map $(0,0,1)^T$ to \vec{k} , $\vec{k} = C(0,0,1)^T$. This implies:

$$\left(\begin{array}{c} k_x \\ k_y \\ k_z \end{array}\right) = \left(\begin{array}{c} a_x \\ a_y \\ a_z \end{array}\right)$$

There are two coordinate systems in which I want to consider rotation of θ radians about \vec{k} . Consider an arbitrary point on the sphere whose C frame coordinates and the base frame coordinates are:

$$\left(\begin{array}{c} x_c \\ y_c \\ z_c \end{array}\right), \quad \left(\begin{array}{c} x_b \\ y_b \\ z_b \end{array}\right) = C \left(\begin{array}{c} x_c \\ y_c \\ z_c \end{array}\right), \left(\begin{array}{c} x_c \\ y_c \\ z_c \end{array}\right) = C^T \left(\begin{array}{c} x_b \\ y_b \\ z_b \end{array}\right)$$

respectively before rotation. They are:

$$\begin{pmatrix} c\theta & -s\theta & 0 \\ s\theta & c\theta & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x_c \\ y_c \\ z_c \end{pmatrix} \text{ and } C \begin{pmatrix} c\theta & -s\theta & 0 \\ s\theta & c\theta & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} x_c \\ y_c \\ z_c \end{pmatrix}$$

after rotation respectively. Hence in terms of Base coordinates:

$$\begin{pmatrix} x_b \\ y_b \\ z_b \end{pmatrix} \longrightarrow C \begin{pmatrix} c\theta & -s\theta & 0 \\ s\theta & c\theta & 0 \\ 0 & 0 & 1 \end{pmatrix} C^T \begin{pmatrix} x_b \\ y_b \\ z_b \end{pmatrix}$$

Let

$$C = \left(\begin{array}{ccc} n_x & o_x & a_x \\ n_y & o_y & a_y \\ n_z & o_z & a_z \end{array}\right)$$

The above product

$$C\begin{pmatrix} c\theta & -s\theta & 0 \\ s\theta & c\theta & 0 \\ 0 & 0 & 1 \end{pmatrix}C^{T} = \begin{pmatrix} n_{x} & o_{x} & a_{x} \\ n_{y} & o_{y} & a_{y} \\ n_{z} & o_{z} & a_{z} \end{pmatrix} \begin{pmatrix} c\theta & -s\theta & 0 \\ s\theta & c\theta & 0 \\ 0 & 0 & 1 \end{pmatrix} \begin{pmatrix} n_{x} & o_{x} & a_{x} \\ n_{y} & o_{y} & a_{y} \\ n_{z} & o_{z} & a_{z} \end{pmatrix}$$
$$= \begin{pmatrix} n_{x}c\theta + o_{x}s\theta & -n_{x}s\theta + o_{x}c\theta & a_{x} \\ n_{y}c\theta + o_{y}s\theta & -n_{y}s\theta + o_{y}c\theta & a_{y} \\ n_{z}c\theta + o_{z}s\theta & -n_{z}s\theta + o_{z}c\theta & a_{z} \end{pmatrix} \begin{pmatrix} n_{x} & n_{y} & n_{z} \\ o_{x} & o_{y} & o_{z} \\ a_{x} & a_{y} & a_{z} \end{pmatrix}$$

The elements of this matrix product can be expressed as:

$$\begin{array}{lll} \operatorname{element}(1,1) & : & n_x^2c\theta + n_xo_xs\theta - n_xo_xs\theta + o_x^2c\theta + a_x^2 & = (n_x^2 + 0_x^2)c\theta + a_x^2 \\ \operatorname{element}(2,1) & : & n_yn_xc\theta + n_xo_ys\theta - n_yo_xs\theta + o_x0_yc\theta + a_xa_y & = (n_xn_y + o_xo_y)c\theta + (n_xo_y - n_yo_x)s\theta + a_xa_y \\ \operatorname{element}(3,1) & : & n_zn_xc\theta + n_xo_zs\theta - n_zo_xs\theta + o_zo_xc\theta + a_xa_z & = (n_zn_x + o_xo_z)c\theta + (n_xo_z - n_zo_x)s\theta + a_xa_z \\ \operatorname{element}(1,2) & : & (n_xn_y + o_xo_y)c\theta + (o_xn_y - n_xo_y)s\theta + a_xa_y \\ \operatorname{element}(2,2) & : & (n_y^2 + o_y^2)c\theta + a_y^2 \\ \operatorname{element}(3,2) & : & (n_zn_y + o_zo_y)c\theta + (n_yo_z - n_zo_y)s\theta + a_ya_z \end{array}$$

Recall that for i = x, y, z we have $n_i^2 + o_i^2 + a_i^2 = 1$, and that $n_i n_j + o_i o_j + a_i a_j = 0$ for $i \neq j$, $(CC^T = I)$.

Furthermore, $\vec{n} \times \vec{o} = \vec{a}$.

$$\begin{pmatrix} k_x \\ k_y \\ k_z \end{pmatrix} = \begin{pmatrix} a_x \\ a_y \\ a_z \end{pmatrix} = \begin{pmatrix} 0 & -n_z & n_y \\ n_z & 0 & -n_x \\ -n_y & n_x & 0 \end{pmatrix} \begin{pmatrix} o_x \\ o_y \\ o_z \end{pmatrix}$$

The very complicated matrix above may be written more simply by eliminating $n'_i s, o'_i s$. In particular, the (1,1) entry can be written as:

$$k_x^2 + (1 - k_x^2)c\theta = c\theta + k_x^2(1 - c\theta)$$

Continuing in this manner we can express every 3×3 rotation in the form:

$$\begin{pmatrix} k_x^2(1-\cos\theta)+\cos\theta & -k_z\sin\theta+k_xk_y(1-\cos\theta) & k_y\sin\theta+k_xk_z(1-\cos\theta) \\ k_z\sin\theta+k_xk_y(1-\cos\theta) & k_y^2(1-\cos\theta)+\cos\theta & -k_x\sin\theta+k_yk_z(1-\cos\theta) \\ -k_y\sin\theta+k_xk_z(1-\cos\theta) & k_x\sin\theta+k_yk_z(1-\cos\theta) & k_z^2(1-\cos\theta)+\cos\theta \end{pmatrix}$$

If we are given:

$$\left(\begin{array}{ccc}
\bar{n}_x & \bar{o}_x & \bar{a}_x \\
\bar{n}_y & \bar{o}_y & \bar{a}_y \\
\bar{n}_z & \bar{o}_z & \bar{a}_z
\end{array}\right)$$

we can solve for $(k_x, k_y, k_z)^T$ and θ . Take the trace of both sides:

$$\bar{n}_x + \bar{o}_y + \bar{a}_z = k_x^2 (1 - \cos \theta) + \cos \theta + k_y^2 (1 - \cos \theta) + \cos \theta + k_z^2 (1 - \cos \theta) + \cos \theta$$

$$= (1 - \cos \theta) + 3\cos \theta$$

$$= 1 + 2\cos \theta$$

$$\implies \cos \theta = \frac{\bar{n}_x + \bar{o}_y + \bar{a}_z - 1}{2}$$

unique solution if $0 \le \theta \le \pi$. Solve for θ . Then solve for k_x, k_y, k_z by looking at the off-diagonal terms.

$$k_z \sin \theta = \frac{\bar{n}_y - \bar{o}_x}{2}$$

$$k_y \sin \theta = \frac{\bar{a}_x - \bar{n}_z}{2}$$

$$k_x \sin \theta = \frac{\bar{o}_z - \bar{a}_y}{2}$$